

PHYS4150 — PLASMA PHYSICS
LECTURE 2 - PLASMA PROPERTIES: DENSITY AND
TEMPERATURE

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Plasma properties: Density and Temperature.

1 PLASMA PROPERTIES

COMPOSITION: ions and electrons

NUMBER DENSITY: ions and electrons in laboratory plasmas $\sim 10^8 - 10^{14} \text{cm}^{-3}$

TEMPERATURE: measured in electron Volts (eV), $1 \text{eV} = 11,600 \text{K}$

$$\begin{aligned} e \cdot U &= \frac{m}{2} v^2 = k_B T = e \cdot 1 \text{V} \\ 1 \text{eV} &= 1.602 \cdot 10^{-19} \text{C} \cdot \text{J/C} \\ 1 \text{eV} &= 1.602 \cdot 10^{-19} \text{J} \end{aligned}$$

DISTANCE SCALE: Debye length λ_D

TIME SCALE: plasma frequency $\omega_p = 2\pi f_p$

VELOCITY SCALE: thermal velocity $v_{th} = \sqrt{\frac{8kT}{\pi m}}$

2 REVIEW: THERMODYNAMICS

Let us starting with a review of some important thermodynamical principles.

2.1 *First law of thermodynamics*

The *First Law of Thermodynamics* states that the change of the internal energy U is given by the sum of the work δW and heat δQ exchanged with the environment:

$$dU = \delta W + \delta Q \quad (1)$$

U is an extensive state function and a thermodynamical potential

Note the use of δ instead of d . This indicates that the amount of exchanged heat

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and work does depend on how the thermodynamical process is performed, and thus, δW and δQ are not *exact differentials*. In contrast, the change of the interior energy depends only on the initial and final state and is therefore an *exact differential*.

2.2 Second law of thermodynamics

The *Second Law of Thermodynamics* is closely related to the *entropy*, which is defined as the reversibly exchanged heat at constant temperature T

S is an extensive state function, while T is an intensive state function

$$dS = \frac{\delta Q}{T}. \quad (2)$$

The second law says now that for a closed system at equilibrium the entropy does not change, i.e.

$$dS = 0. \quad (3)$$

At a given temperature the amount of irreversibly exchanged heat is always smaller than the amount of reversibly exchanged heat, and thus

$$\delta Q_{irr} < \delta Q_{rev} = T dS. \quad (4)$$

For a closed system at equilibrium the entropy takes its maximum value S_{max} , while for an irreversible process $dS > 0$.

2.3 Ideal gas

In an ideal gas the particles are assumed to undergo only elastic collisions. In this case the equation of state is

$$pV = Nk_B T, \quad (5)$$

where p , V , and N are the pressure, volume, and particle number of the gas. The Boltzmann constant k_B

$$k_B = 1.308 \cdot 10^{-23} \text{J/K} = 8.617 \cdot 10^{-5} \text{eV}. \quad (6)$$

relates the average kinetic energy of the gas with the temperature. For an ideal gas the average (translational) energy is

$$\frac{1}{2} m \langle v^2 \rangle = \frac{3}{2} k_B T. \quad (7)$$

3 DENSITY

SOLID As an example let us consider aluminum which has a density of $\rho_{Al} = 3 \cdot 10^3 \text{kg/m}^3$ and an atomic mass of $m_{Al} = 27u$. We now want to find the number of aluminum atoms per unit volume: $u = 1.66 \cdot 10^{-27} \text{kg}$ is the atomic mass unit.

$$n_{Al} = \frac{\rho_{Al}}{m_{Al}u} = \frac{3 \cdot 10^3 \text{kg/m}^3}{27 \cdot 1.66 \cdot 10^{-27} \text{kg}} = 6.8 \cdot 10^{28} \text{m}^{-3}. \quad (8)$$

AIR At standard pressure one mol of air has a volume of $22.4 \text{l} = 22.4 \cdot 10^{-3} \text{m}^3$. One mol are $6 \cdot 10^{23}$ particles, and thus

$$n_{air} = \frac{6 \cdot 10^{23}}{22.4 \cdot 10^{-3} \text{m}^3} = 2.7 \cdot 10^{25} \text{m}^{-3}. \quad (9)$$

	$n[\text{m}^{-3}]$	kT [eV]
Solar wind @ Earth	5	50
ionosphere	$10^5 - 10^6$	0.02
Solar corona	10^6	100
tokamak	10^{14}	10^4
laser-produced	10^{20}	100
glow discharge	$10^8 - 10^{10}$	2

4 TEMPERATURE

Let us have a closer look at the velocity distribution $f(\mathbf{v})$ of a gas and how it relates to its temperature. Because the gas motion is isotropic, $f(\mathbf{v})$ can only be a function of \mathbf{v}^2 . On the other hand, the components of $f(\mathbf{v})$ must be independent, which implies that

$$f(\mathbf{v}^2) = f(v_x^2 + v_y^2 + v_z^2) = f(v_x^2)f(v_y^2)f(v_z^2). \quad (10)$$

The only function that fulfills Eq. (10) is

$$f(\mathbf{v}^2) = c \cdot e^{-a\mathbf{v}^2}. \quad (11)$$

To find the constant c we require that the components of f are normalized, i.e. $\int f_i(v_i) dv = 1$, which is only possible if $a < 0$, and

$$1 = c \int e^{-av^2} dv = c \sqrt{\frac{\pi}{a}}. \quad (12)$$

To obtain the constant a we note that in a gas at equilibrium the energy per degree of freedom is $\frac{1}{2}k_B T$, and therefore

$$k_B T = m \langle v_i^2 \rangle = m \int v_i^2 f(v_i) dv_i = m \sqrt{\frac{\pi}{a}} \int \exp\{-av_i^2\} v_i^2 dv_i. \quad (13)$$

$$dv_i = \frac{1}{2\sqrt{a}} \frac{dx}{\sqrt{x}}$$

Replacing the argument of the exponential by $x = av_i^2$ we get

$$k_B T = \frac{m}{\sqrt{\pi a}} \int_0^\infty e^{-x} \sqrt{x} dx = \frac{m}{\sqrt{\pi a}} \Gamma\left(\frac{3}{2}\right), \quad (14)$$

where the *Gamma function* $\Gamma(x)$ is defined as

$$\Gamma(z) = \int_0^\infty e^{-x} x^{z-1} dx \quad (15)$$

$$\Gamma(z+1) = \Gamma(z) \cdot z \quad (16)$$

$$\Gamma(1) = 1 \quad (17)$$

$$\Gamma\left(\frac{1}{2}\right) = \sqrt{\pi} \quad (18)$$

From this follows that $\Gamma(3/2) = \frac{\sqrt{\pi}}{2}$ and

$$f(v) = \sqrt{\frac{m}{2\pi k_B T}} \exp\left\{-\frac{mv^2}{2k_B T}\right\}$$

$$f(\mathbf{v}) = \left\{\frac{m}{2\pi k_B T}\right\}^{3/2} \exp\left\{-\frac{m\mathbf{v}^2}{2k_B T}\right\}$$